

# A PHS Tour in Audio/Acoustics

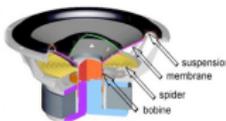
## *Applications in Virtual Reality, Control and Health*

Thomas Hélie, CNRS

S3AM Team

Laboratory of Sciences and Technologies of Music and Sound  
IRCAM – CNRS – Sorbonne Université  
Paris, France

Spring School on Theory and Applications of Port-Hamiltonian Systems  
31 March 2019 – 5 April 2019  
Fraueninsel (Chiemsee)



## Context and motivation

- 1 **Model input-output multi-physics systems** for sound and musical applications:
  - Phenomena: mechanical, acoustic, electronic, magnetic, etc
  - Realism: nonlinearities, non ideal dissipations, etc
- 2 **Satisfy fundamental physical properties**:
  - causality, stability, passivity and more precisely ...
  - the **power balance** structured into **conservative**/**dissipative**/**source** parts
  - other natural invariants and symmetries (if any)
- 3 **Simulate such systems and preserve these properties** in the discrete time domain (*+accuracy+sound quality/Shannon-Nyquist principle*)
- 4 **Design code generators** from netlists for real-time applications
- 5 **Design correctors and controllers** to reach target behaviours

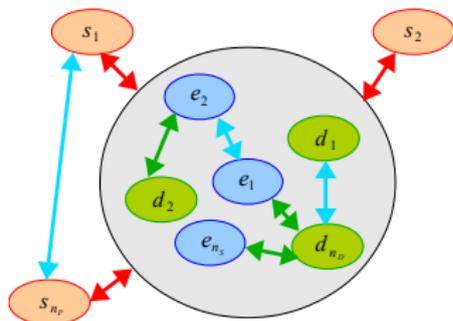
# Outline

- 1 Context
- 2 **Framework:** basics, recalls and tools
- 3 **Analog electronics** and electro-acoustics
- 4 **Mechanics:** nonlinear damped vibrations
- 5 **Control** applications in acoustics
- 6 **Voice:** a minimal model to analyze self-oscillations
- 7 Conclusion

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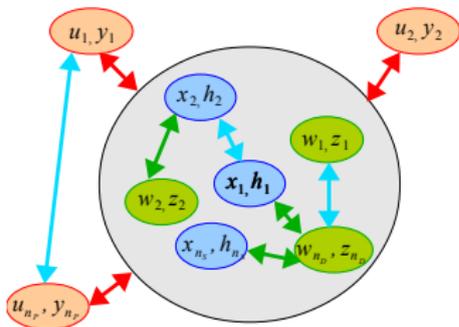
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  - Modelling: Component-based approach & Port-Hamiltonian Systems
  - Power-balanced numerical method
  - Tool: the PyPHS Python library [Falaize]
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## A physical system is made of...



- **Energy-storing components:** (energy)  
 $E = \sum_{n=1}^N e_n \geq 0$
- **Dissipative components:** (dissipated power)  
 $Q = \sum_{m=1}^M d_m \geq 0$
- **External sources:** (external power)  
 $P_{\text{ext}} = \sum_{p=1}^P s_p$
- **Conservative connections** (power balance)  
 $\frac{dE}{dt} = P_{\text{ext}} - Q$

# A physical system is made of...



- Energy-storing components:** (energy)  
 $E = H(\mathbf{x}) = \sum_{n=1}^N H_n(x_n) \geq 0$
- Dissipative components:** (dissipated power)  
 $Q = \mathbf{z}(\mathbf{w})^T \mathbf{w} = \sum_{m=1}^M z_m(w_m) w_m \geq 0$   
 (effort  $\times$  flow : force  $\times$  velocity, voltage  $\times$  current, etc)
- External sources:** (external power)  
 $P_{\text{ext}} = \mathbf{u}^T \mathbf{y} = \sum_{p=1}^P u_p y_p$
- Conservative connections** (power balance)  
 $\nabla H(\mathbf{x})^T \frac{d\mathbf{x}}{dt} + \mathbf{z}(\mathbf{w})^T \mathbf{w} - \mathbf{u}^T \mathbf{y} = 0$

## Port-Hamiltonian Formulation

## Power balance

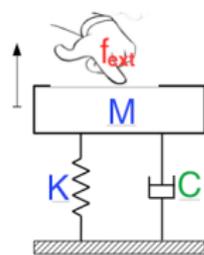
$$\underbrace{\begin{pmatrix} \frac{d\mathbf{x}}{dt} \\ \mathbf{w} \\ -\mathbf{y} \end{pmatrix}}_B = S \cdot \underbrace{\begin{pmatrix} \nabla H(\mathbf{x}) \\ \mathbf{z}(\mathbf{w}) \\ \mathbf{u} \end{pmatrix}}_A$$

$$\begin{aligned} 0 &= A^T B \\ &= A^T S A \end{aligned}$$

if  $S = -S^T$

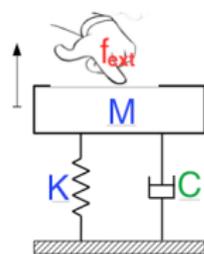
A: components constitutive laws & external actions,  
 S: interconnections between flows and efforts.

## Example: damped mechanical oscillator



$\mathbf{x} = \begin{pmatrix} p \\ q \end{pmatrix}$	<i>momentum elongation</i>	$H(\mathbf{x}) = \frac{1}{2} \mathbf{x}^T \begin{pmatrix} M^{-1} & 0 \\ 0 & K \end{pmatrix} \mathbf{x}$	<i>kinetic potential</i>
$\frac{d\mathbf{x}}{dt} = \begin{pmatrix} \dot{p} \\ \dot{q} \end{pmatrix}$	<i>inertia force spring velocity</i>	$\nabla H(\mathbf{x}) = \begin{pmatrix} p/M \\ K q \end{pmatrix}$	<i>mass velocity spring force</i>
$\mathbf{w} = v_d$	<i>dashpot velocity</i>	$\mathbf{z}(\mathbf{w}) = C v_d$	<i>damping force</i>
$\mathbf{y} = v_{\text{ext}}$	<i>external velocity</i>	$\mathbf{u} = f_{\text{ext}}$	<i>external force</i>

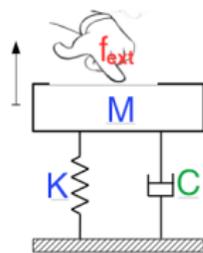
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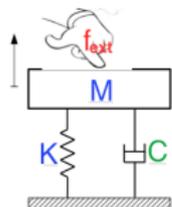


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$$\begin{aligned}
 \dot{p} &= f_{\text{ext}} - K q - C v_d \\
 \dot{q} &= p/M \\
 v_d &= p/M \\
 v_{\text{ext}} &= p/M
 \end{aligned}$$

$$\begin{pmatrix} \frac{d\mathbf{x}}{dt} \\ \mathbf{w} \\ -\mathbf{y} \end{pmatrix} = \underbrace{\begin{pmatrix} 0 & -1 & -1 & +1 \\ +1 & 0 & 0 & 0 \\ +1 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}}_{S = -S^T} \cdot \begin{pmatrix} \nabla H(\mathbf{x}) \\ \mathbf{z}(\mathbf{w}) \\ \mathbf{u} \end{pmatrix}$$

## Some variations



$$\begin{pmatrix} F_M \\ v_K \\ v_C \\ -v_{\text{ext}} \end{pmatrix} = \begin{pmatrix} 0 & -1 & -1 & +1 \\ +1 & 0 & 0 & 0 \\ +1 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} \cdot \begin{pmatrix} v_M \\ F_K \\ F_C \\ F_{\text{ext}} \end{pmatrix}$$

Hamiltonian systems (conservative, autonomous)

$$\begin{pmatrix} F_M \\ v_K \\ \cdot \\ \cdot \end{pmatrix} = \begin{pmatrix} 0 & -1 & \cdot & \cdot \\ +1 & 0 & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \end{pmatrix} \cdot \begin{pmatrix} v_M \\ F_K \\ \cdot \\ \cdot \end{pmatrix}$$

"Mass+Damper+Excitation" (spring removed)

$$\begin{pmatrix} F_M \\ \cdot \\ v_C \\ -v_{\text{ext}} \end{pmatrix} = \begin{pmatrix} 0 & \cdot & -1 & +1 \\ \cdot & \cdot & \cdot & \cdot \\ +1 & \cdot & 0 & 0 \\ -1 & \cdot & 0 & 0 \end{pmatrix} \cdot \begin{pmatrix} v_M \\ \cdot \\ F_C \\ F_{\text{ext}} \end{pmatrix}$$

"Mass+Excitation"

$$\begin{pmatrix} F_M \\ \cdot \\ \cdot \\ -v_{\text{ext}} \end{pmatrix} = \begin{pmatrix} 0 & \cdot & \cdot & +1 \\ \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot \\ -1 & \cdot & \cdot & 0 \end{pmatrix} \cdot \begin{pmatrix} v_M \\ \cdot \\ \cdot \\ F_{\text{ext}} \end{pmatrix}$$

# Formulations

## 1. Differential-Algebraic

$$\begin{pmatrix} \frac{dx}{dt} \\ \mathbf{w} \\ -\mathbf{y} \end{pmatrix} = S \cdot \begin{pmatrix} \nabla H(\mathbf{x}) \\ \mathbf{z}(\mathbf{w}) \\ \mathbf{u} \end{pmatrix}, \quad S = -S^T$$

## 2. Differential (1 → 2) by solving algebraic part $\mathbf{w} = W(\nabla H(\mathbf{x}), \mathbf{u})$

$$\begin{cases} \frac{dx}{dt} = (J - R) \nabla H(\mathbf{x}) + \mathbf{G}u, & J = -J^T, \quad R = R^T \geq 0 \\ -\mathbf{y} = -\mathbf{G}^T \nabla H(\mathbf{x}) + D\mathbf{u}, & D = -D^T \end{cases}$$

"Mass-Spring-Damper":  $H(\mathbf{x}) = \frac{x_1^2}{2M} + \frac{Kx_2^2}{2}$ ,  $\mathbf{z}(\mathbf{w}) = C\mathbf{w}$

$$S = \begin{bmatrix} 0 & -1 & -1 & +1 \\ +1 & 0 & 0 & 0 \\ +1 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{bmatrix}, \quad J = \begin{bmatrix} 0 & -1 \\ +1 & 0 \end{bmatrix}, \quad R = \begin{bmatrix} C & 0 \\ 0 & 0 \end{bmatrix}, \quad \mathbf{G} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}, \quad D = 0$$

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Classical numerical schemes for  $\frac{dx}{dt} = f(x)$ :

- efficiently approximate  $\frac{d}{dt}$  and exploit  $f$
- *a posteriori* analysis of stability

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A discrete power-balanced method (PHS)

Exploit differentiation chain rule

$$\frac{dE}{dt} = \sum_n \frac{\partial H}{\partial x_n} \frac{dx_n}{dt} \approx \sum_n \underbrace{\frac{H_n(x_n[k+1]) - H_n(x_n[k])}{x_n[k+1] - x_n[k]}}_{\nabla^d H(x[k], \delta x[k])_n} \underbrace{\frac{x_n[k+1] - x_n[k]}{\delta t}}_{(\delta x[k]/\delta t)_n} = \frac{E[k+1] - E[k]}{\delta t}$$

Jointly substitute  $\dot{x} \rightarrow \delta x/\delta t$  and  $\nabla H(x) \rightarrow \nabla^d H(x, \delta x)$ :

$$\underbrace{\begin{pmatrix} \frac{\delta x}{\delta t} \\ \mathbf{w} \\ -\mathbf{y} \end{pmatrix}}_B = S \underbrace{\begin{pmatrix} \nabla^d H(x, \delta x) \\ \mathbf{z}(\mathbf{w}) \\ \mathbf{u} \end{pmatrix}}_A$$

Simulation : solve  $(\delta x, w)$  at each time step  $k$  (e.g. Newton-Raphson algo.)

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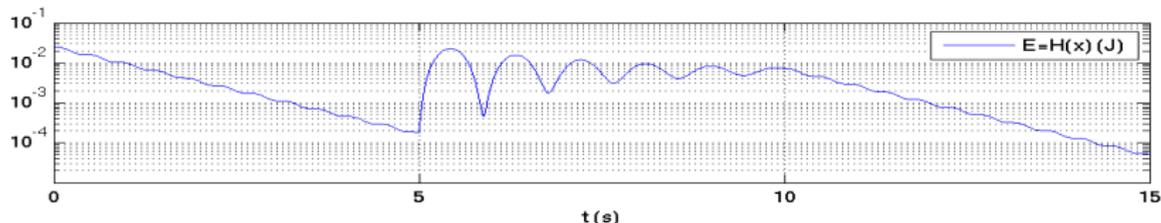
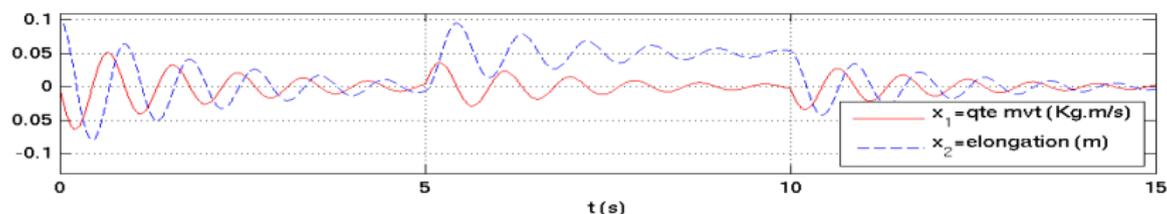
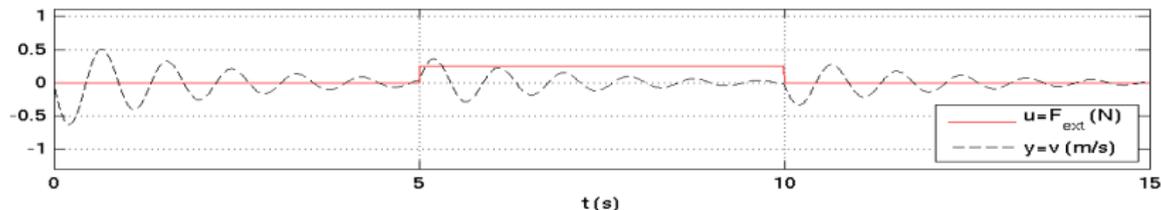
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Simulation : solve  $(\delta x, \mathbf{w})$  at each time step  $k$  (e.g. Newton-Raphson algo.)

- Skew-symmetry of  $S$  preserved  $\Rightarrow 0 = A^T S A = A^T B = \delta E/\delta t + \mathbf{z}(\mathbf{w})^T \mathbf{w} - \mathbf{u}^T \mathbf{y}$
- For **linear systems**,  $\nabla^d H(x, \delta x) = \nabla(x + \delta x/2)$  restores the **mid-point scheme**.
- Method also applies to nonlinear components and non separate Hamiltonian
- Also available: Power-balanced Runge-Kutta scheme (non iterative algo.) [Lopes et al., LHMNLC'2015]

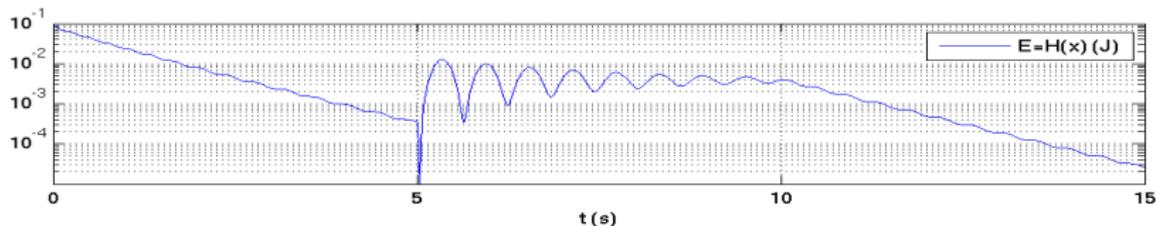
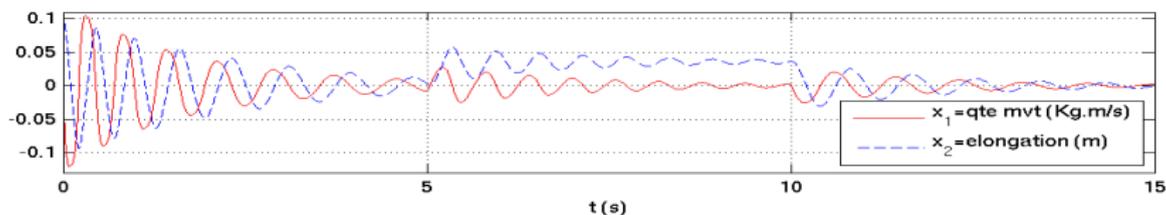
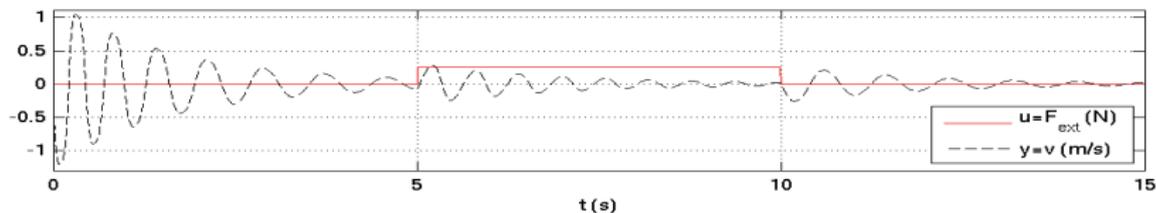
# Simulation 1: mass-spring-damper

- Parameters:  $M=100$  g,  $K=5$  N/m,  $C=0.1$  N.s/m et  $\delta t=5$  ms
- Initial conditions:  $x_0 = [mv_0=0, \ell_0=10 \text{ cm}]^T$
- Excitation:  $F_{\text{ext}}(t) = F_{\text{max}} \mathbf{1}_{[5\text{s},10\text{s}]}(t)$  with  $F_{\text{max}} = K\ell_0/2 = 0.25$  N



## Simulation 2: idem with a hardening spring

- **Potential energy:**  $H_2^{\text{NL}}(x_2) = K L^2 [\cosh(x_2/L) - 1]$  ( $\sim k x_2^2/2$ )
- **Physical law:**  $F_2 = (H_2^{\text{NL}})'(x_2) = K L \sinh(x_2/L)$  ( $\sim K x_2$ )
- **Reference elongation:**  $L = \ell_0/4 = 25 \text{ mm}$



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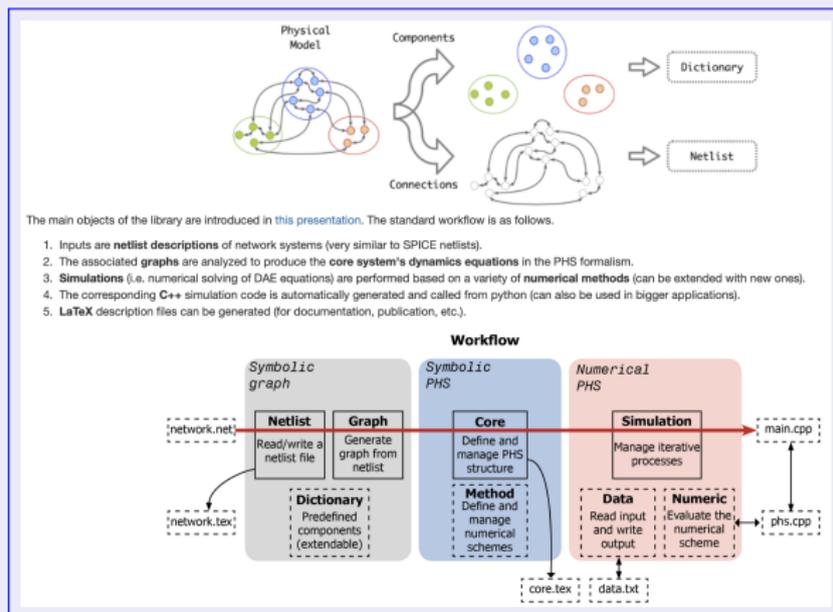
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<https://pyphs.github.io/pyphs/>

2012-16 : First version

[Falaize, PhD]

2016-- : Opensource library with periodic releases [Falaize & contributors]



→ a short presentation (pdf file)

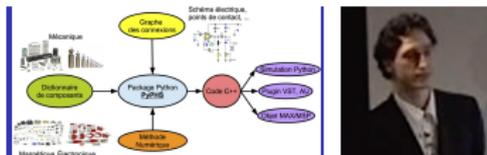
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  - Guitar Pedal and Electric Piano
  - Ondes Martenot
  - Operational Amplifier
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[Falaize, PhD'16]  
[Najnudel et al., AES'18]  
[Muller et al., DAFx'19]

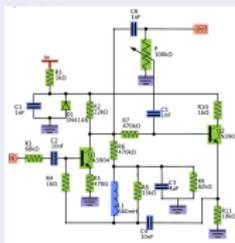
# PhD, 2016: Antoine Falaize

Passive modelling, simulation, code generation and correction of audio multi-physical systems



## Two examples

Wah pedal (CryBaby): netlist → **PyPHS** → LaTeX eq. & C code



Components	Number
Storage	7 linear
Dissipative	18 (5 NL, 2 modulated)
Ports	3 (IN, OUT, battery)

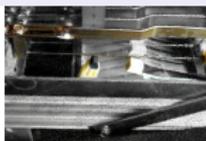
Audio PlugIn:

Sound 1a: dry

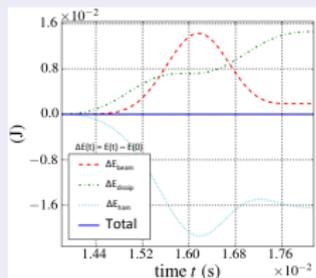
Sound 1b: wah

## A simplified Fender-Rhodes Piano

Sound 2



Components	Hammer	1 beam	Pickup/RC-circuit
Storage	2 NL	2M lin.	2 lin. (+ NL connection)
Dissipative	1 NL	M lin.	1 lin.
Ports	2	1	1



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## Ondes Martenot

(created by Maurice Martenot in 1928)

Controls



Circuit



*Diffuseurs*



→ Video 3 [Thomas Bloch, improvisation, 2010]

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Context/Problem

(Musée de la Musique, Philharmonie de Paris)

Technological obsolescence of a musical instrument:

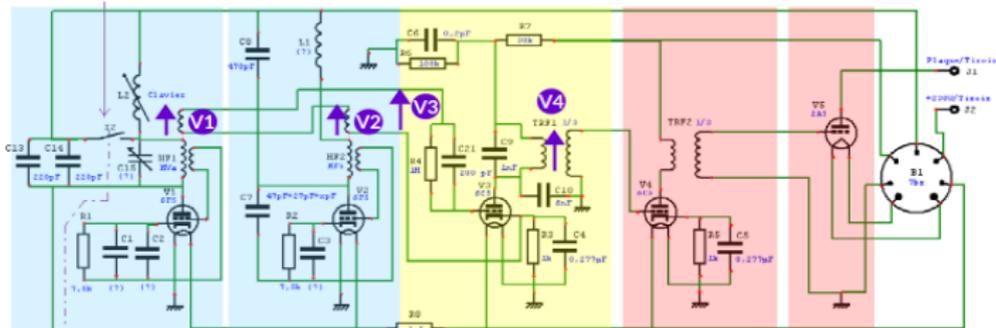
**70/281 remaining instruments** (*handmade*), **1200 pieces** (*Varèse, Maessian, etc*)

Objective

(Collegium Musicae-Sorbonne Université)

Real-time simulation of the circuit based on physics → PHS approach

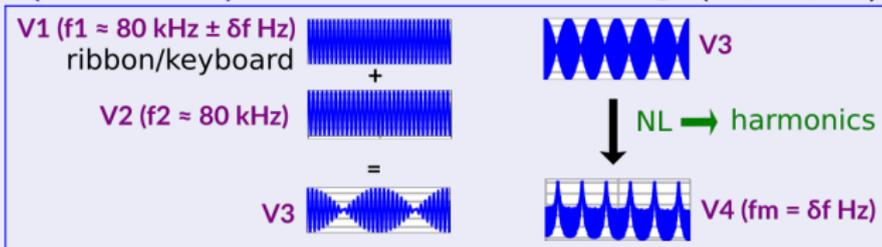
# Ondes Martenot: 5 stages circuit



var. osc. fixed osc. demodulator preamp. power amp.

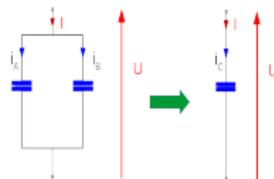
## Specificities: heterodyne oscillators (1930's)

- 2 High frequencies ( $\approx 80\text{kHz} \pm \delta f$ )  $\rightarrow$  demodulator  $\rightarrow$  audio range ( $\delta f, 2\delta f, \dots$ )



- Vacuum tubes:  $w = [\text{grid and plate currents}]^T$ ,  $z(w) = \text{associated voltages}$  (passive parametric model [Cohen'12])
- **Pb:** ribbon-controlled oscillator involving **time-varying capacitors in parallel**

## Ondes Martenot: capacitors in parallel



Problem:

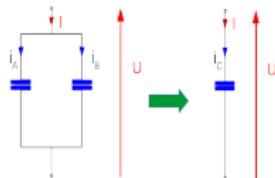
Capacitors	$(n = A, B)$
State ( <i>charge</i> ):	$q_n$
Energy :	$H_n(q_n)$
Flux ( <i>current</i> ):	$i_n = dq_n/dt$
Effort ( <i>voltage</i> ):	$v_n = H'_n(q_n)$

$$v_C = v_A = v_B \quad \&$$

$$\begin{bmatrix} i_A \\ i_B \\ -v_C \end{bmatrix} = \begin{bmatrix} \text{not} \\ \text{realizable} \end{bmatrix} \begin{bmatrix} v_A = H'_A(q_A) \\ v_B = H'_B(q_B) \\ i_C \end{bmatrix}$$

→ Build the equivalent component  $C = A//B$

# Ondes Martenot: capacitors in parallel



Problem:

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→ Build the equivalent component  $C = A // B$

Hyp:  $q_n \mapsto v_n = H'_n(q_n)$  bijective (increasing law)

Find the total energy  $H_C(q_C)$  for the total charge  $q_C = q_A + q_B$

- ① Charge as a function of the voltage  $v_n = v_C$ :  $q_n = [H'_n]^{-1}(v) := Q_n(v_C)$
- ② Total charge (idem):  $q_C = [Q_A + Q_B](v_C) := Q_C(v_C)$
- ③ Total energy function:  $H_C(q_C) = \sum_{n=A,B} H_n \circ Q_n \circ Q_C^{-1}(q_C)$

Also available if  $H_n$  depends on additional states (ribbon position  $\ell$ )

Power-balanced simulation

with  $H(q, \ell) = q^2 / (2C_{\text{Martenot}}(\ell))$

→ video 4 (sound=circuit output voltage, without the *diffuseurs*)

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- 6 **Voice:** a minimal model to analyze self-oscillations
- 7 Conclusion

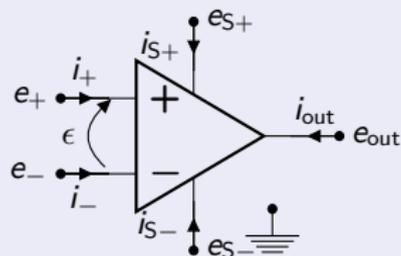
[Falaize, PhD'16]

[Najnudel et al., AES'18]

[Muller et al., DAFx'19]

## Idealized component

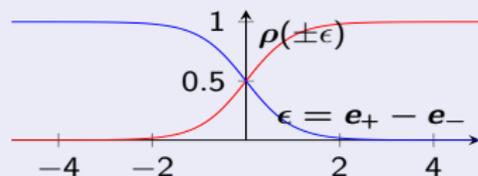
## • 5 ports



## • Algebraic conservative law

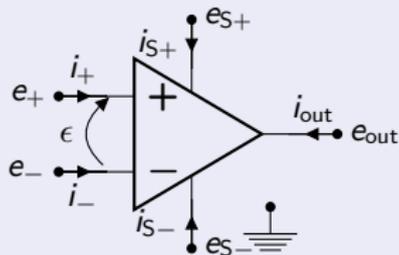
$$\underbrace{\begin{bmatrix} i_+ \\ i_- \\ i_{S+} \\ i_{S-} \\ e_{out} \end{bmatrix}}_{Z_{OPA}(w)} = \underbrace{\begin{bmatrix} \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & -\rho(+\epsilon) & \cdot \\ \cdot & \cdot & \cdot & -\rho(-\epsilon) & \cdot \\ \cdot & \cdot & \rho(\epsilon) & \rho(-\epsilon) & \cdot \end{bmatrix}}_{J(W_{OPA}) = -J(W_{OPA})^T} \underbrace{\begin{bmatrix} e_+ \\ e_- \\ e_{S+} \\ e_{S-} \\ i_{out} \end{bmatrix}}_{W_{OPA}}$$

## • Modulation factor



## Idealized component

- 5 ports

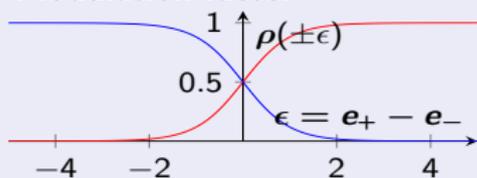


- Algebraic conservative law

$$\underbrace{\begin{bmatrix} i_+ \\ i_- \\ i_{S+} \\ i_{S-} \\ e_{out} \end{bmatrix}}_{ZOPA(w)} = \underbrace{\begin{bmatrix} \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \end{bmatrix}}_{J(WOPA) = -J(WOPA)^T} \underbrace{\begin{bmatrix} e_+ \\ e_- \\ e_{S+} \\ e_{S-} \\ i_{out} \end{bmatrix}}_{WOPA}$$

$\rho(+\epsilon)$      $\rho(-\epsilon)$      $\rho(\epsilon)$      $\rho(-\epsilon)$

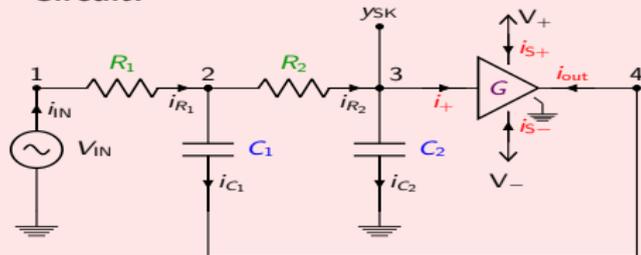
- Modulation factor



## Typical analog filters

## (Sallen-Key)

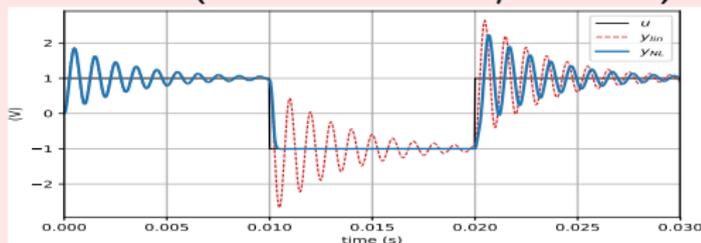
- Circuit:



- Nonlinear PHS:

$$\begin{bmatrix} \dot{x} \\ w \\ WOPA \\ -y \end{bmatrix} = S \begin{bmatrix} \nabla H(x) \\ z(w) \\ ZOPA(WOPA) \\ u \end{bmatrix}$$

- Sounds 5 (simulations: **linear** / **nonlinear**)



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[Hélie, Matignon, 2015]

[Hélie, Roze '18]

[Raibeau, Roze '18]

## 1. Theoretical issues

Given a linear conservative mechanical system,

- find **damping models** that preserve the eigen modes (with eigen structure)
- design **nonlinear** damping in such a class
- provide a **power balanced formulation** that is preserved in **simulations**

## 2. Application in musical acoustics

Build **physical models** to produce:

- a **variety of beam sounds** (glockenspiel, xylophone, marimba, etc)
- **morphed sounds** through some extrapolations based on **physical grounds**  
(e.g. *meta-materials with damping depending on the magnitude*)

## Damping models for $M\ddot{q} + C\dot{q} + Kq = f$ (finite-dimensional case)

### Conservative problem ( $C=0$ )

- $\ddot{q} + (M^{-1}K)q = M^{-1}f$
- Eigen-modes  $e_j$ :  $(M^{-1}K)e_j = \omega_j^2 e_j$  ( $\omega_j$ : angular freq.)

### Damping that preserves eigen-modes ?

- Choose  $M^{-1}C$  as a **non-negative polynomial** of matrix  $M^{-1}K$
- *Caughey class (1960)*:  $C = c_0M + c_1K + c_2KM^{-1}K + \dots$

## Damping models for $M\ddot{q} + C\dot{q} + Kq = f$ (finite-dimensional case)

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→ *Caughey class (1960)*:  $C = c_0M + c_1K + c_2KM^{-1}K + \dots$

### Eigen-modes with nonlinearly-damped dynamics ?

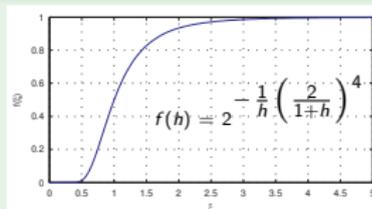
- Make  $c_n$  depend on the dynamics

### Ex.: damping as a function of energy $H(x)$

(state  $x = [q, p = M\dot{q}]^T$ )

$c_n(x) = \kappa_n(H(x)) \in [c_n^-, c_n^+]$  with  $c_n^- \geq 0$

- Increasing:  $\kappa_n(h) = c_n^- + (c_n^+ - c_n^-)f\left(\frac{h}{h_0}\right)$
- Decreasing:  $\kappa_n(h) = c_n^+ - (c_n^+ - c_n^-)f\left(\frac{h}{h_0}\right)$



## Application: the Euler-Bernoulli beam

### 1. Pinned beam excited by a distributed force

(H1) Euler-Bernoulli kinematics: straight cross-section after deformation

(H2) linear approximation for the conservative problem

(H3) viscous and structural dampings: only  $c_0, c_1 \geq 0$

### 2. Dimensionless model

( $w$ : deflection,  $t \geq 0$ ,  $0 \leq \ell \leq 1$ )

- **PDE:**  $\underbrace{\partial_t^2 w}_{\mathcal{M} \equiv Id} + \underbrace{(c_0 + c_1 \partial_\ell^4)}_c \partial_t w + \underbrace{\partial_\ell^4 w}_{\mathcal{K}} = f_{\text{ext}}$

- **Boundaries**  $\ell \in \{0, 1\}$ : fixed extremities ( $w=0$ ), no momentum ( $\partial_\ell^2 w=0$ )

- **Energy:**  $E = \int_0^1 \left( \frac{(\partial_\ell^2 w)^2}{2} + \frac{(\partial_t w)^2}{2} \right) d\ell$

### 3. Modal decomposition: $e_m(\ell) = \sqrt{2} \sin(m\pi\ell)$

( $1 \leq m \leq n$ )

PHS:  $\partial_t x = (J - R)\nabla H(x) + Gu$  with  $J = \begin{bmatrix} 0_{n \times n} & I_n \\ -I_n & 0_{n \times n} \end{bmatrix}$ ,  $R = \begin{bmatrix} 0_{n \times n} & 0_{n \times n} \\ 0_{n \times n} & C \end{bmatrix}$

$y = G^T \nabla H(x)$   $G^T = [0_{n \times n}, I_n]$

with  $H(x = [q; p = M\dot{q}]) = \frac{1}{2} p^T M^{-1} p + \frac{1}{2} q^T K q$

and  $q = [q_1, \dots, q_n]^T$ ,  $u = [u_1, \dots, u_n]^T$ ,  $y = [y_1, \dots, y_n]^T$

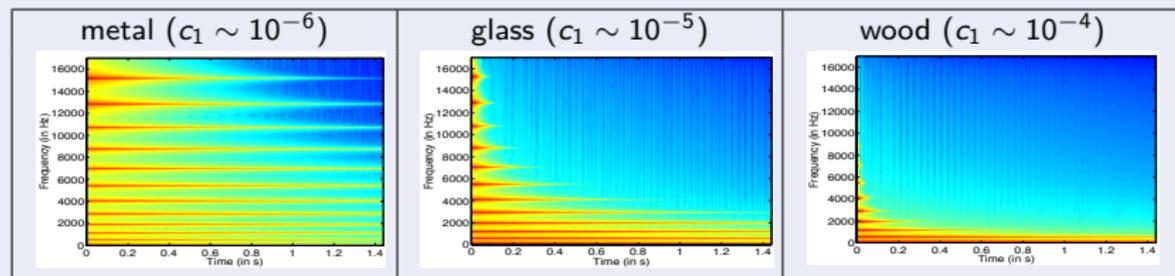
(projections of  $w$ ,  $f_{\text{ext}}$ ,  $v_{\text{ext}}$ )

where  $M = I_n$ ,  $K = \pi^4 \text{diag}(1, \dots, n)^4$  and  $C = c_0 I_n + c_1 K$ .

# Damping and simulation parameters

Examples of spectrograms for standard linear dampings:

$$c_0 \sim 10^{-2}$$



Nonlinear damping (from metal to wood):

$$C(x) = c_0(x)I + c_1(x)K \text{ with}$$
$$c_n(x) = \beta_n(H(x)) \in [c_n^-, c_n^+]$$

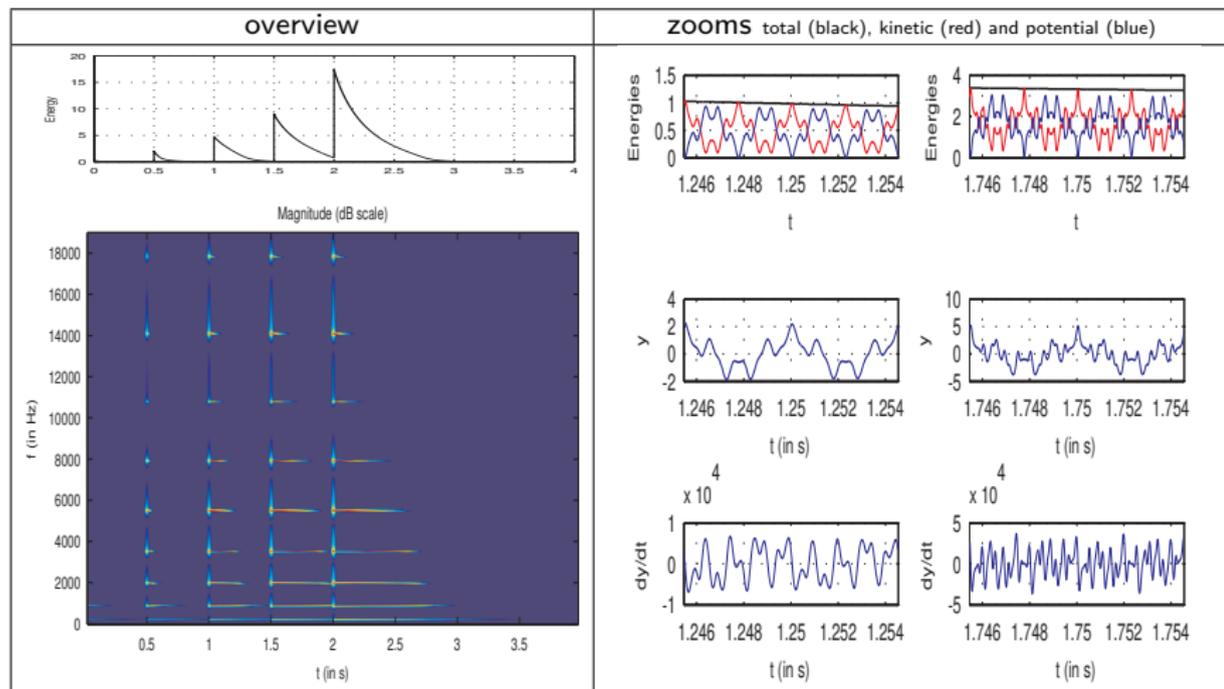
metal	$c_0^- = 0.02$	$c_1^- = 10^{-6}$
wood	$c_0^+ = 0.04$	$c_1^+ = 10^{-4}$

Numerical method preserving the power balance (discrete gradient)

- force distributed close to  $z = 0$ :  $u = [1, \dots, 1]^T f$
- listened signal: acceleration  $[1, \dots, 1]\dot{y}$
- $n = 9$  modes and time step s.t.  $f_1 = 220$  Hz to  $f_9 \approx n^2 f_1 = 17820$  Hz

Results:  $H(x) \ll 1 \rightarrow$  wood,  $H(x) \gg 1 \rightarrow$  metal

force: 5 piecewise constant pulses (0.1ms) with increasing magnitude  $\rightarrow$  Sound 6



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[Hélie, Matignon, 2015]

[Hélie, Roze '18]

[Raibeau, Roze '18]

## 1. Dimensionless model

$$(x, y) \in \Omega = (0, 1)^2, t \geq 0$$

- PDE:  $\partial_t^2 w + \alpha \partial_t w + \Delta \Delta w - \epsilon \left( \int_{\Omega} |\nabla w|^2 dS \right) \Delta w = f$
- Zero IC & pinned BC ( $w = 0$  on  $\partial\Omega$ ,  $\partial_x^2 w = 0$  if  $x \in \{0, 1\}$ ,  $\partial_y^2 w = 0$  if  $y \in \{0, 1\}$ )

$$\text{Energy: } E = \int_{\Omega} \frac{(\partial_t^2 w)^2}{2} dS + \int_{\Omega} \frac{(\Delta w)^2}{2} dS + \epsilon \left( \int_{\Omega} \frac{(\nabla w)^2}{2} dS \right)^2$$

2. Modal decomposition (if  $\epsilon = 0$ ):  $e_{kl}(x, y) = 2 \sin(k\pi x) \sin(l\pi y)$  ( $1 \leq k, l \leq n$ )

$$\text{PHS: } \partial_t x = (J - R) \nabla H(x) + Gu \quad \text{with } J = \begin{bmatrix} 0_{n \times n} & I_n \\ -I_n & 0_{n \times n} \end{bmatrix}, R = \begin{bmatrix} 0_{n \times n} & 0_{n \times n} \\ 0_{n \times n} & C \end{bmatrix}$$

$$y = G^T \nabla H(x) \quad G^T = [0_{n \times n}, I_n]$$

→ idem that the previous problem with

$$K = \pi^4 \text{diag}(1, \dots, (k^2 + l^2)^2, \dots, (n^2 + n^2)^2) \quad \text{and } C = \alpha I_n$$

## 1. Dimensionless model

$$(x, y) \in \Omega = (0, 1)^2, t \geq 0$$

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$$K = \pi^4 \text{diag}(1, \dots, (k^2 + l^2)^2, \dots, (n^2 + n^2)^2) \text{ and } C = \alpha I_n$$

## Nonlinear case: $\epsilon > 0$

- Modal decomposition available:  $\epsilon \left( \int_{\Omega} |\nabla w|^2 dS \right) \Delta w$  colinear to  $\Delta w$
- Replace quadratic  $H$  by  $H(q, p) = \frac{1}{2} p^T p + \frac{1}{2} q^T K q + \epsilon \left( \frac{1}{2} q^T K^{1/2} q \right)^2$   
 → pitch effect on sounds → video example 7 with sound and spectrum

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[Hélie, Roze '18]

**[Raibeau, Roze '18]**

## Nonlinear Kirchhoff-Carrier string model: PDE → FEM

$$\partial_t^2 w + \alpha \partial_t w - \left( 1 + \varepsilon \int_0^1 (\partial_\ell w)^2 d\ell \right) \partial_\ell^2 w = f$$

$$\xrightarrow{\text{FEM}} \mathbf{M}\ddot{\mathbf{W}} + \mathbf{C}\dot{\mathbf{W}} + (1 + \beta \mathbf{W}^T \mathbf{K} \mathbf{W}) \mathbf{K} \mathbf{W} = \mathbf{F} \quad \text{with } \mathbf{W} = [w_1 w_2 \dots w_n]^T.$$

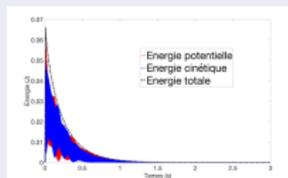
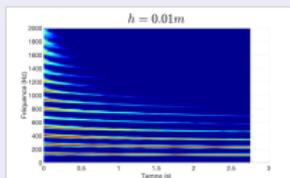
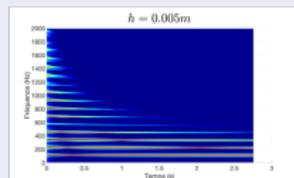
$$\text{Energy: } H(Q = \mathbf{W}, P = \mathbf{M}\dot{\mathbf{W}}) = \frac{1}{2} \mathbf{P}^T \mathbf{M}^{-1} \mathbf{P} + \left( 1 + \frac{\beta}{2} \mathbf{Q}^T \mathbf{K} \mathbf{Q} \right) \frac{1}{2} \mathbf{Q}^T \mathbf{K} \mathbf{Q}.$$

## FEM → PHS (Finite Elements Method in the PHS formalism)

$$\underbrace{\begin{bmatrix} \dot{Q} \\ \dot{P} \end{bmatrix}}_{\dot{\mathbf{X}}} = \left( \underbrace{\begin{bmatrix} 0 & I \\ -I & 0 \end{bmatrix}}_J - \underbrace{\begin{bmatrix} 0 & 0 \\ 0 & C \end{bmatrix}}_R \right) \underbrace{\begin{bmatrix} (1 + \beta \mathbf{Q}^T \mathbf{K} \mathbf{Q}) \mathbf{K} \mathbf{Q} \\ \mathbf{M}^{-1} \mathbf{P} \end{bmatrix}}_{\nabla H} + \underbrace{\begin{bmatrix} 0 \\ I \end{bmatrix}}_G \underbrace{F}_u$$

## Energy-preserving simulations of a nonlinear dynamics

→ sounds 8



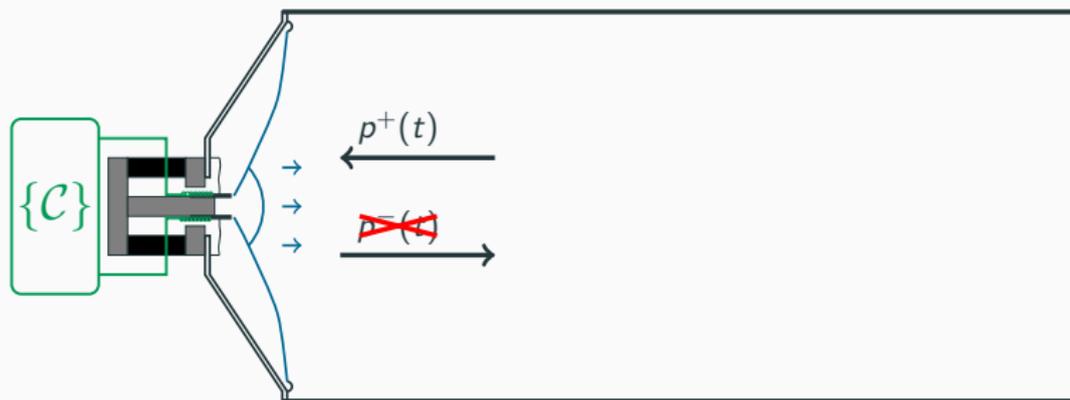
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  - Passive Finite-Time Control of loudspeakers [Lebrun, Wijnand et al., Nodycon'19]
  - (zoom) Robotised testbed for brass instruments [Lopes, PhD'16]
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## Motivation: sound absorption



Here:  $\{C\}$

(a) Finite-time control ( $\neq$  linear stabilization)

(b) Preserve controller passivity

Combination of (a) and (b): not straightforward.

## Plane wave propagation in a tube

- 1D plane waves propagation

$$p(z, t) = p^+(t - z/c_0) + p^-(t + z/c_0)$$

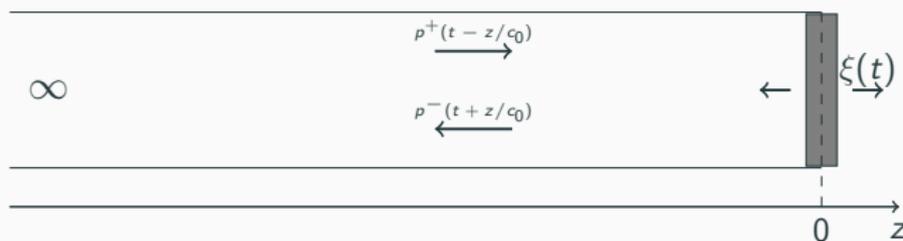
$$v(z, t) = \frac{p^+(t - z/c_0) - p^-(t + z/c_0)}{\rho_0 c_0}$$

- Boundary condition at  $z = 0$ : rigid piston

$$p^-(0, t) = p^+(0, t) - \rho_0 c_0 \dot{\xi}(t)$$

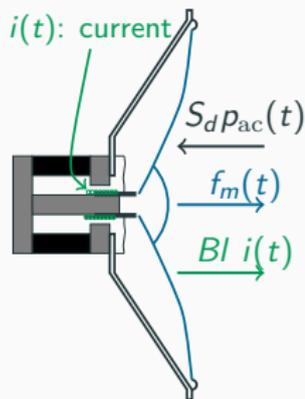
→ Impedance matching condition:

$$\dot{\xi}^*(t) = \frac{p(0, t)}{\rho_0 c_0}$$



## Current-driven loudspeaker

$$S_d p_{ac}(t) = \underbrace{M_m \ddot{\xi}(t) + R_m \dot{\xi}(t) + K_m \xi(t)}_{f_m(t)} + Bl i(t)$$



- Laplace force
- Force due to the mechanical subsystem
- Force due to the acoustic pressure

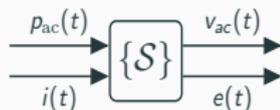
## Power-balanced formulation

$$\text{Stored energy: } \mathcal{H}(x) = K_m \frac{\xi^2}{2} + \frac{p^2}{2M_m} \geq 0$$

$$\underbrace{\begin{bmatrix} \dot{\xi} \\ \dot{p} \end{bmatrix}}_{\dot{x}} = \underbrace{\begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix}}_{J=-J^T} - \underbrace{\begin{bmatrix} 0 & 0 \\ 0 & R_m \end{bmatrix}}_{R \geq 0} \underbrace{\begin{bmatrix} K_m \xi \\ p/M_m \end{bmatrix}}_{\nabla \mathcal{H}(x)} + \underbrace{\begin{bmatrix} 0 & 0 \\ -Bl & S_d \end{bmatrix}}_G \underbrace{\begin{bmatrix} i \\ p_{ac} \end{bmatrix}}_u$$

$$\underbrace{\begin{bmatrix} e \\ v_{ac} \end{bmatrix}}_y = \underbrace{\begin{bmatrix} 0 & -Bl \\ 0 & S_d \end{bmatrix}}_{G^T} \underbrace{\begin{bmatrix} K_m \xi \\ p/M_m \end{bmatrix}}_{\nabla \mathcal{H}(x)}$$

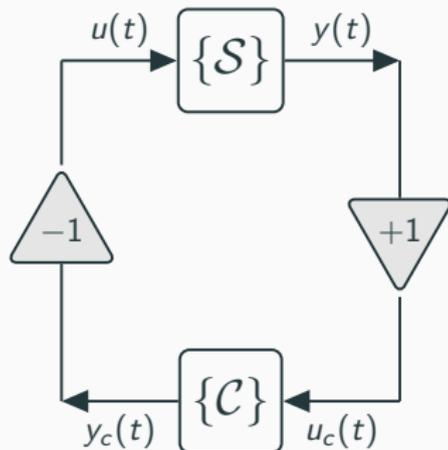
**Port-Hamiltonian System**  
[VanderShaft2014]



## Ingredient 1: passive control

## (PHS approach)

## A passive controller ?



- $\{S\}$  is a port-Hamiltonian system (PHS):

$$\{S\} : \begin{cases} \dot{x}(t) = (\mathbf{J} - \mathbf{R}) \nabla \mathcal{H}(x) + \mathbf{G}u(t) \\ y(t) = \mathbf{G}^T \nabla \mathcal{H}(x) \end{cases}$$

- **Passivity:**

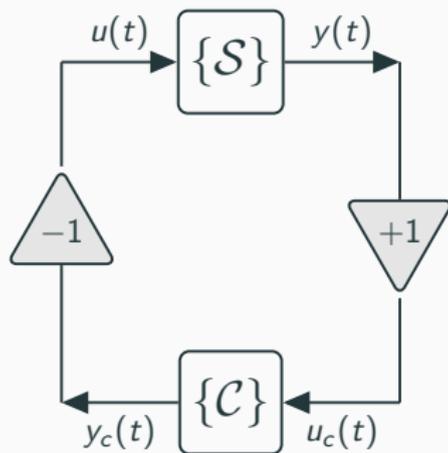
$$\underbrace{\nabla \mathcal{H}(x)^T \frac{dx}{dt}}_{\text{stored power}} \leq \underbrace{y^T u}_{\text{external power}}$$

- Same property imposed to  $\{C\}$  :

$$\{C\} : \begin{cases} \dot{x}_c(t) = (\mathbf{J}_c - \mathbf{R}_c) \nabla \mathcal{H}_c(x_c) + \mathbf{G}_c u(t) \\ y_c(t) = \mathbf{G}_c^T \nabla \mathcal{H}_c(x_c) \end{cases}$$

## Ingredient 1: passive control

## A passive controller ?



Interconnection of  $\{S\}$  &  $\{C\}$  is a PHS:

$$\begin{bmatrix} \dot{x} \\ \dot{x}_c \end{bmatrix} = \begin{bmatrix} \mathbf{J} - \mathbf{R} & -\mathbf{G}\mathbf{G}_c^T \\ \mathbf{G}_c\mathbf{G}^T & \mathbf{J}_c - \mathbf{R}_c \end{bmatrix} \nabla \mathcal{H}_{s+c}(x, x_c)$$

where  $\mathcal{H}_{s+c}(x, x_c) = \mathcal{H}_s(x) + \mathcal{H}_c(x_c)$ .

## Example of control design

1. Define a total energy  $\mathcal{H}_{s+c}(x, x_c)$  that has an equilibrium at a **target  $x^*$** .

$$\mathcal{H}_{s+c}(x, x_c) = (x - x^*)^2.$$

2. Deduce the energy of the controller

$$\mathcal{H}_c(x_c) = \mathcal{H}_{s+c}(x, x_c) - \mathcal{H}(x) \geq 0$$

## Ingredient 2: finite-time control

### Finite-time control of a double integrator [Bernuau2015]

Consider system

$$\begin{cases} \dot{z}_1 = z_2 \\ \dot{z}_2 = v \end{cases}$$

Using the control law for  $0 < \alpha < 1$

$$v = -k_1 [z_1]^{2-\alpha} - k_2 [z_2]^\alpha$$

with  $[x]^\alpha \triangleq \text{sign}(x)|x|^\alpha$ ,  $k_1 > 0$ ,  $k_2 > 0$ , the origin is **globally finite-time stable**.

- **Finite-time control**  $\rightarrow$  **nonlinear control**
- **Reaches target in finite-time** ( $\neq$  asymptotic convergence)
- × **Controller not passive**

## Passive finite-time control

### Application to the loudspeaker system

- Target closed-loop energy:

$$\mathcal{H}_{s+c}(\xi, p) = M_m k_1 \frac{2-\alpha}{2} |\xi - \xi^*|^{\frac{2}{2-\alpha}} + \frac{M_m k_2}{R_m} \frac{1}{\alpha+1} \left| \frac{p - p^*}{M_m} \right|^{\alpha+1}$$

→ Minimum at  $(\xi, p) = (\xi^*, p^*)$

→  $\mathcal{H}_{s+c} - H \neq$  positive definite

## Passive finite-time control

### Application to the loudspeaker system

- Target closed-loop energy:

$$\mathcal{H}_{s+c}(\xi, p) = M_m k_1 \frac{2-\alpha}{2} |\xi - \xi^*|^{\frac{2}{2-\alpha}} + \frac{M_m k_2}{R_m} \frac{1}{\alpha+1} \left| \frac{p - p^*}{M_m} \right|^{\alpha+1} + \frac{\beta}{2} (\xi - \xi^*)^2 + \frac{\gamma}{2M_m} (p - p^*)^2.$$

→ Minimum at  $(\xi, p) = (\xi^*, p^*)$

→  $\beta > K_m$  and  $\gamma > R_m$  for **controller passivity**

- Energy of the controller:  $\mathcal{H}_{s+c} - H_s = \mathcal{H}_c \geq 0$

$$\mathcal{H}_c(\xi, p) = M_m k_1 \frac{2-\alpha}{2} |\xi - \xi^*|^{\frac{2}{2-\alpha}} + \frac{M_m k_2}{R_m} \frac{1}{\alpha+1} \left| \frac{p - p^*}{M_m} \right|^{\alpha+1} + \frac{\beta}{2} (\xi - \xi^*)^2 + \frac{\gamma}{2M_m} (p - p^*)^2 - \frac{1}{2M_m} p^2 - \frac{K_m}{2} \xi^2 \geq 0$$

## Passive finite-time control

### Application to the loudspeaker system

- Controller  $\{\mathcal{C}\}$  :

$$\dot{x} = \left( \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} - \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix} \right) \nabla \mathcal{H}_c(x) + \begin{bmatrix} \frac{1}{Bl} \\ \frac{R_m}{Bl} \end{bmatrix} u_c$$

$$y_c = \begin{bmatrix} \frac{1}{Bl} & \frac{R_m}{Bl} \end{bmatrix} \nabla \mathcal{H}(x_c),$$

$\rightarrow \beta > K_m$  and  $\gamma > R_m$  for **controller passivity**

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## Numerical results

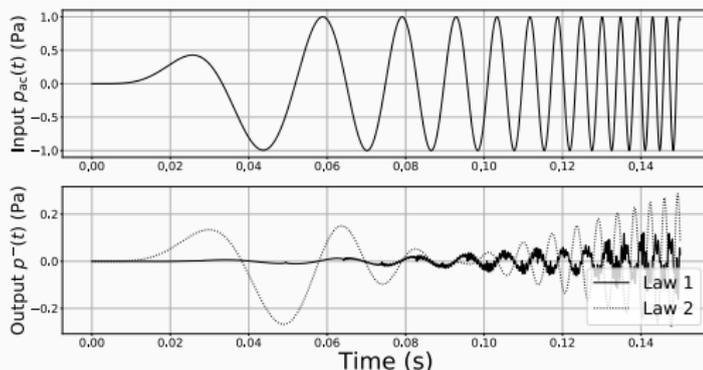
### Simulation configuration

- Power-balanced numerical scheme
- Input: logarithmic chirp  $p_{ac}(t)$
- Output: reflected pressure  $p^-(t)$

### Test cases

- Law 1: proposed passive finite-time ( $\rightarrow$  nonlinear) control
- Law 2: (reference) linear impedance-based control

### Time domain simulation



Absorption  $\geq 75\%$   
in any case

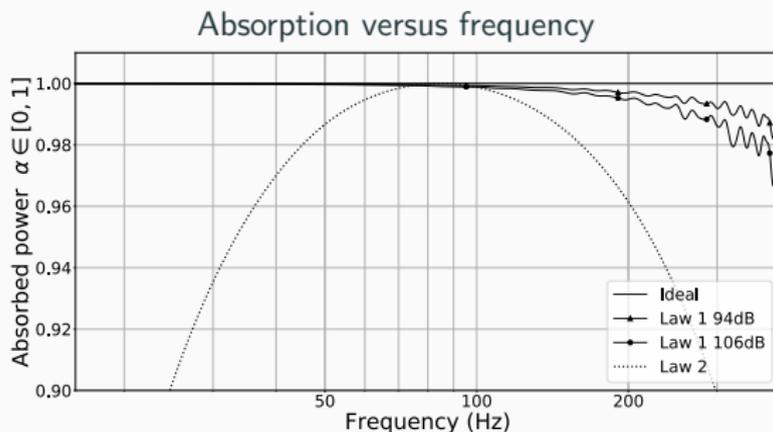
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### Simulation configuration

- Power-balanced numerical scheme
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### Test cases

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## Conclusion & perspectives

### Contribution

- ✓ Construction and simulation of a passive finite-time control law
- ✓ Efficient sound absorption in the low-frequency audio range

### Perspectives

- Stiff problem of finite-time control around the origin (not Lipschitz continuous)
- Hardware implementation: passive delayed interconnection (current work)

# Outline

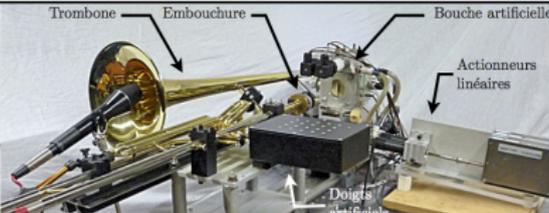
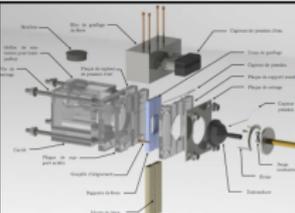
- 1 Context
- 2 **Framework:** basics, recalls and tools
- 3 **Analog electronics** and electro-acoustics
- 4 **Mechanics:** nonlinear damped vibrations
- 5 **Control applications in acoustics**
  - Passive Finite-Time Control of loudspeakers [Lebrun, Wijnand et al., Nodycon'19]
  - (zoom) Robotised testbed for brass instruments [Lopes, PhD'16]
- 6 **Voice:** a minimal model to analyze self-oscillations
- 7 Conclusion

# PhD, June 2016: Nicolas Lopes

## Passive modelling, simulation and experimental study of a robotised artificial mouth playing brass instruments



### 1. PHS/Simu of the complete system: *air jet in a channel with mobile walls, etc*

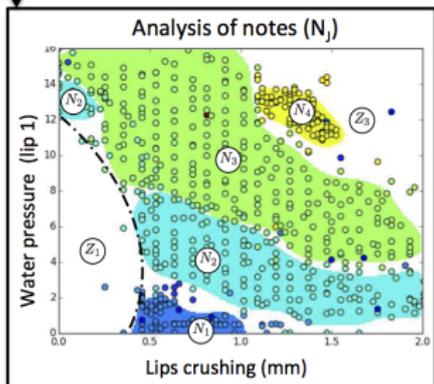


Instrumented mouth, lips and mouthpiece

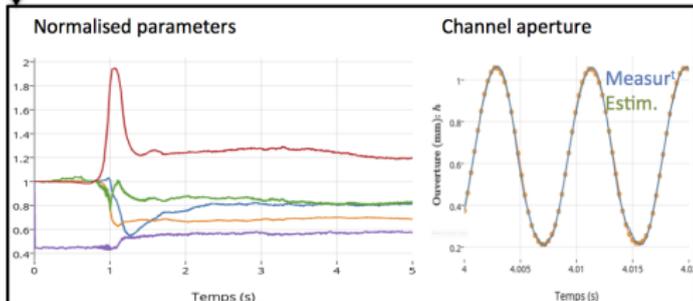
Robotised testbed

HMI (dSpace & Python)

### 2. Automated exploration



### 3. Estimation/Observation: extended Kalman filter on PHS (5 estimated parameters: lip mass, stiffness, etc)



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[Hélie,Silva,Wetzel, 2019]

# Vocal apparatus: Physiology & Physics

Air is forced out of the **lungs** ...

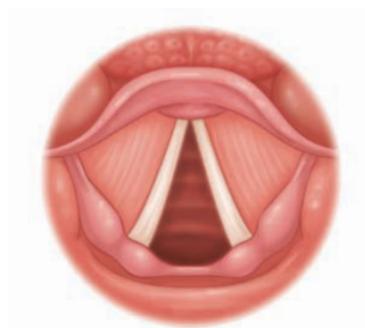
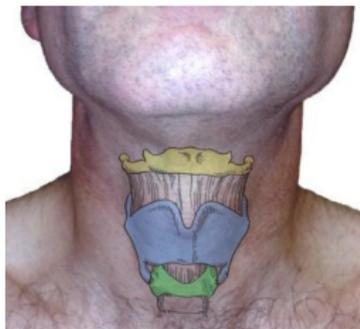
(*active*)

... at the top of the trachea, it goes through the **glottis**,  
(=constriction between the vocal folds)

(*passive*)

... before reaching the **vocal tract** (=pharynx, nose, mouth).

(*passive*)



Source: Seikel et al

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(active)

(passive)

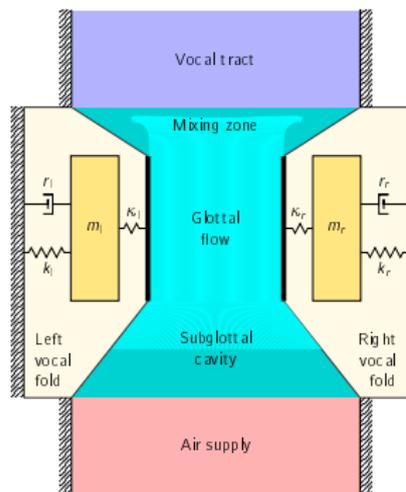
(passive)

**Phonation** due to a **nonlinear** fluid/structure interaction in the larynx

Unstable equilibrium beyond some subglottal pressure threshold:

- Vibration of the folds
- Modulation of the glottal flow
- Coupling with acoustic waves in the vocal tract

Most of state-of-the-art physics-based voice models are **not power-balanced** and **violate passivity**.



# Vocal apparatus: Physiology & Physics

Air is forced out of the **lungs** ...

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... before reaching the **vocal tract** (=pharynx, nose, mouth).

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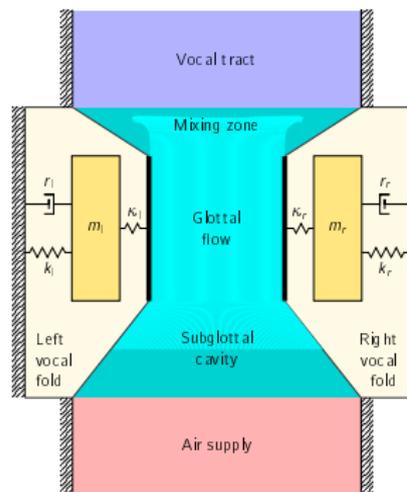
(passive)

(passive)

## Objective

Derive a **minimal model of the voice production** that

- restores a **power balance**
- is structured into components
- enables **power-balanced** time-domain simulations and analyse its bifurcations and oscillation regimes



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[Hélie, Silva, Wetzel, 2019]

# Vocal folds model

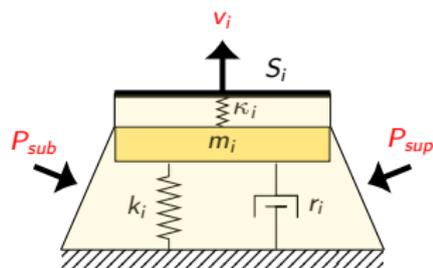
$(i = l, r)$

Each fold is modelled as a spring-mass-damping system to be coupled to the glottal flow through a cover spring.

- Momentum  $\pi_i$  and elongations  $\xi_i$  and  $\eta_i$

$$H = \frac{1}{2} \left( \frac{\pi_i^2}{m_i} + k_i \xi_i^2 + \kappa_i \eta_i^2 \right)$$

- Dissipation  $w_i = \pi_i / m_i$  and  $z_i(w_i) = r_i w_i$
- Flow-controlled at the glottis port ( $v_i$ )
- Effort-controlled at the subglottal ( $P_{sub}$ ) and supraglottal ( $P_{sup}$ ) ports.



$$\begin{bmatrix} \dot{\xi}_i \\ \dot{\pi}_i \\ \dot{\eta}_i \\ w_i \\ -Q_{sub}^i \\ Q_{sup}^i \\ F_i \end{bmatrix} = \begin{bmatrix} -1 & 1 & 1 & -1 & -S_{sub}^i & -S_{sup}^i \\ -1 & & & & & 1 \\ 1 & & & & & \\ S_{sub}^i & & & & & \\ S_{sup}^i & & & & & \\ & & & -1 & & \end{bmatrix} \cdot \begin{bmatrix} \frac{\partial \xi_i}{\partial H} \\ \frac{\partial \pi_i}{\partial H} \\ \frac{\partial \eta_i}{\partial H} \\ z_i(w_i) \\ P_{sub} \\ P_{sup} \\ v_i \end{bmatrix}$$

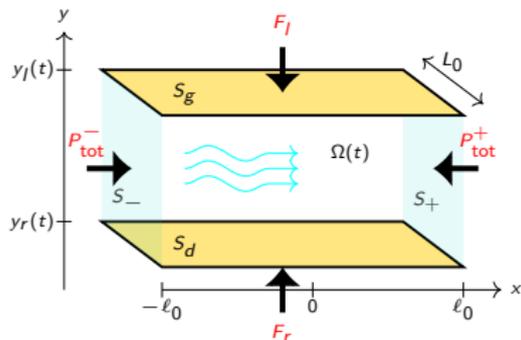
$F_i$  Force of the vocal fold on the glottal flow

$Q_{sup}^i$  Flowrate pushed into the supraglottal cavity

$Q_{sub}^i$  Flowrate pulled from the subglottal cavity

# Glottal flow between two parallel mobile walls

Objective: account for the **transverse velocity** and for **power exchanged on the walls**



Kinematics:

$$\mathbf{v} = \begin{pmatrix} V_x(t) \\ V_y(t) \end{pmatrix} + \frac{V_{exp}(t)}{h(t)} \begin{pmatrix} -x \\ y - y_m \end{pmatrix}$$

$V_x(t)$  mean axial velocity

$V_y(t) = \dot{y}_m$  mean transverse velocity

$V_{exp}(t) = \dot{h}$  transverse expansion velocity

$h(t)$  channel height

## 2D field: Euler equation

$$\nabla \cdot \mathbf{v} = 0, \quad \nabla \times \mathbf{v} = 0 \text{ on } \Omega(t)$$

$$\dot{\mathbf{v}} + \nabla \left( \frac{p}{\rho_0} + \frac{|\mathbf{v}|^2}{2} \right) = 0$$

$$\text{BCs } \mathbf{v}_y(y_r) = \dot{y}_r \quad \mathbf{v}_y(y_l) = \dot{y}_l$$

$$\text{with } h = y_l - y_r \text{ et } y_m = \frac{y_l + y_r}{2}$$

## Reduced order dynamics

$$m(h) \dot{V}_x(t) = L_0 h (P_{tot}^- - P_{tot}^+)$$

$$m(h) \dot{V}_y(t) = F_r - F_l$$

$$m_e(h) \dot{V}_{exp}(t) = L_0 \ell_0 (P_{tot}^- + P_{tot}^+) - \frac{F_r + F_l}{2} - \partial_h H$$

$$\dot{h} = \frac{1}{m_e(h)} \partial_{V_{exp}} H$$

$$H(V_x, V_y, V_{exp}, h) = \frac{1}{2} m(h) (V_x^2 + V_y^2) + \frac{1}{2} m_e(h) V_{exp}^2$$



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[Hélie,Silva,Wetzel, 2019]

# Power-balanced numerical experiments

Parameters values from literature:

**Folds:**  $m_i = 0.2 \text{ g}$ ,  $r_i = 0.05 \text{ N s m}^{-1}$ ,  $\kappa_i = 3k_i$ ,

$S_{\text{sup}}^i = 1.1 \text{ mm}^2$ ,  $S_{\text{sub}}^i: 1.1 \text{ cm}^2$

**Glottis:**  $L_0 = 11 \text{ mm}$ ,  $2\ell_0 = 4 \text{ mm}$

**Vocal tract:** first resonance of /a/

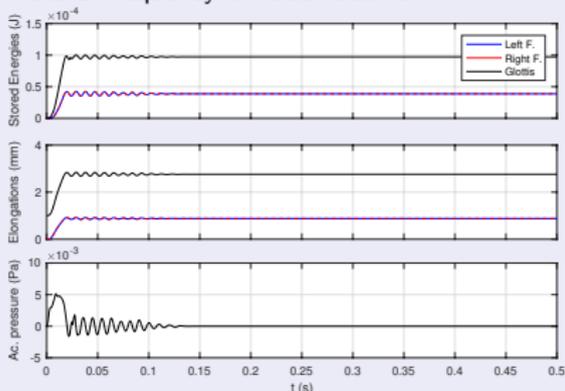
$f_0 = 640 \text{ Hz}$ ,  $q_0 = 2.5$ ,  $Z_0 = 1 \text{ M}\Omega$

**Input**  $P_{\text{sub}}: 0 \nearrow 800 \text{ Pa}$  within 20 ms

Open glottis ( $h_r = 1 \text{ mm}$ )

$k_r = k_l = 100 \text{ N m}^{-1}$

Natural frequency of folds: 112 Hz



No oscillation onset

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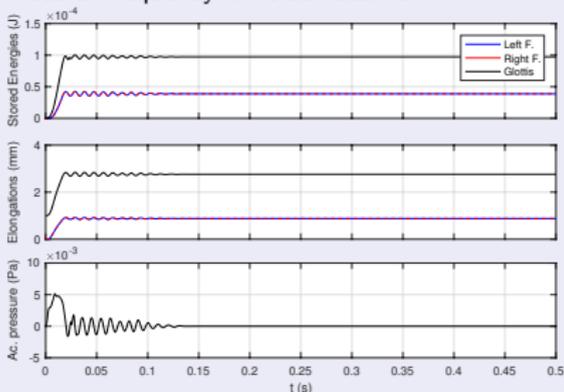
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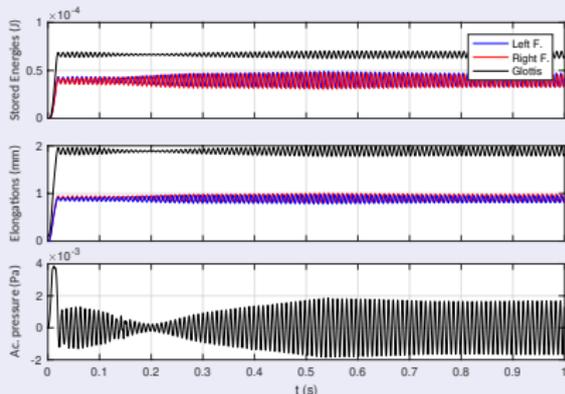
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## Adduction (nearly closed, $h_r = 0.1 \text{ mm}$ )

Slightly detuned vocal folds

$$k_l = 100 \text{ N m}^{-1} \text{ (112 Hz)}$$

$$k_r = 97 \text{ N m}^{-1} \text{ (110 Hz)}$$



Periodic oscillations:

- Oscillation stabilized after some transient
- Synchronized folds vibrations even without contact between folds

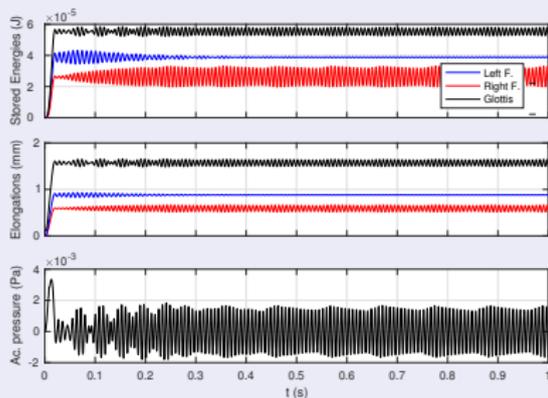
# Power-balanced numerical experiments

## Strong asymmetry

$$k_l = 100 \text{ N m}^{-1} \text{ (112 Hz)}$$

$$k_r = 149 \text{ N m}^{-1} \text{ (137 Hz)}$$

with adduction ( $h_r = 0.1 \text{ mm}$ )



### Quasi-periodic oscillations:

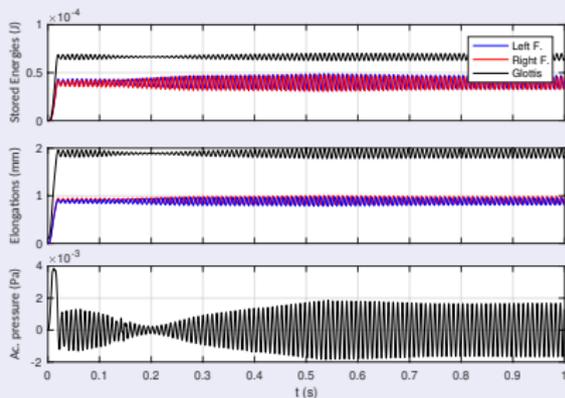
- starting on the left (lax) vocal fold,
- transferred to the right (stiffer) fold for the steady state regime.

## Adduction (nearly closed, $h_r = 0.1 \text{ mm}$ )

Slightly detuned vocal folds

$$k_l = 100 \text{ N m}^{-1} \text{ (112 Hz)}$$

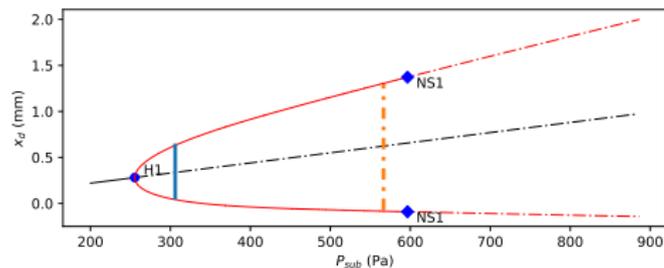
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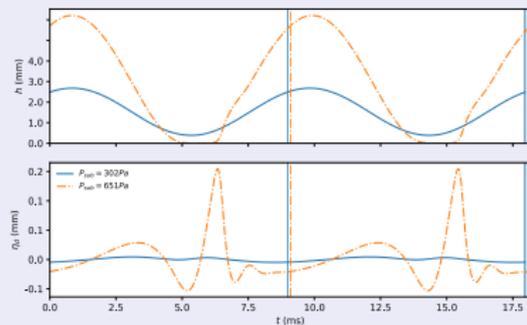
## Coupling PyPHS with PyDSTool (continuation tool)



## Hopf bifurcation

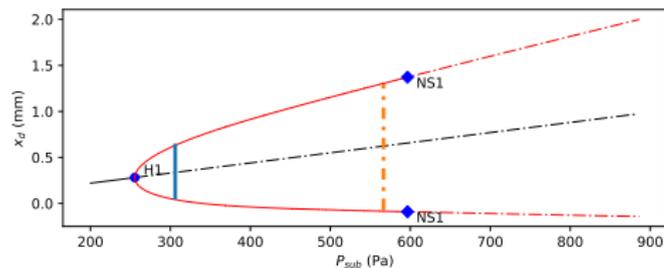
Onset pressure very close to the one obtained from time-domain simulation using PyPHS.

## Beyond the Hopf bifurcation: limit cycles



## Quasi sinusoidal oscillation near threshold

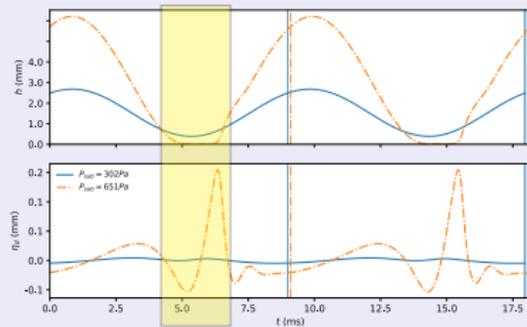
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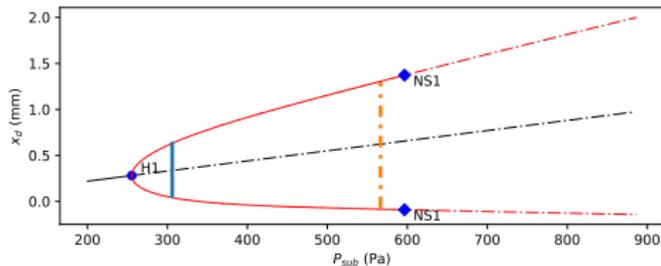
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Force during glottal closure

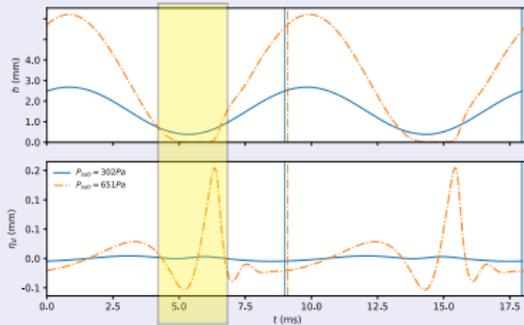
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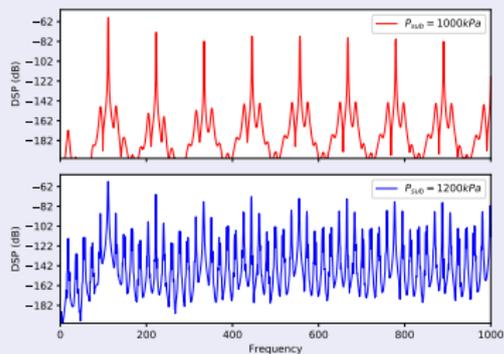
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Force during glottal closure

## Beyond the Neimark-Sacker bifurcation

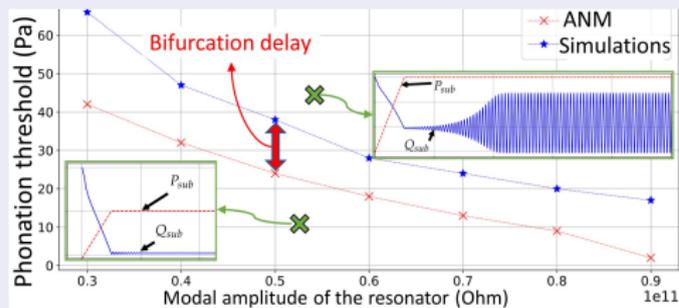


High pressure → raucous sound

# Interpretation of other results on bifurcation analysis

## Oscillation thresholds

- × Continuation of Hopf points (ANM)
- ★ Estimation from time-domain simulations

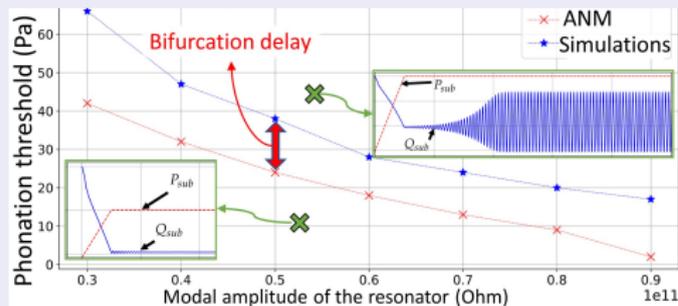


- Explicit bifurcation delay
- Related to dynamic bifurcation phenomenon?

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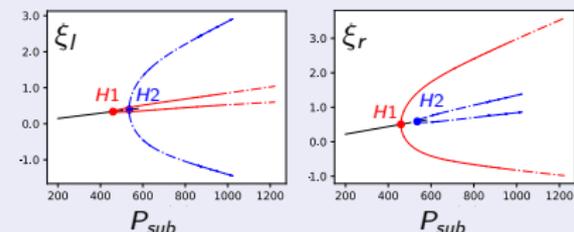
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## Asymmetric vocal folds (L. Forma)

Bifurcation diagrams of  $\xi_l$  and  $\xi_r$

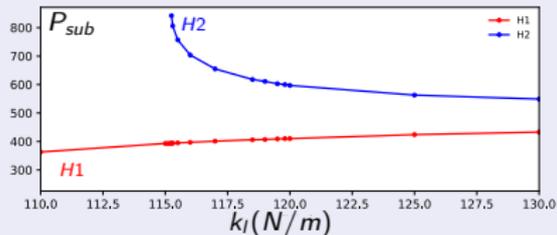
$$k_l = 150N/m$$

$$k_r = 100N/m$$



## Second Hopf bifurcation

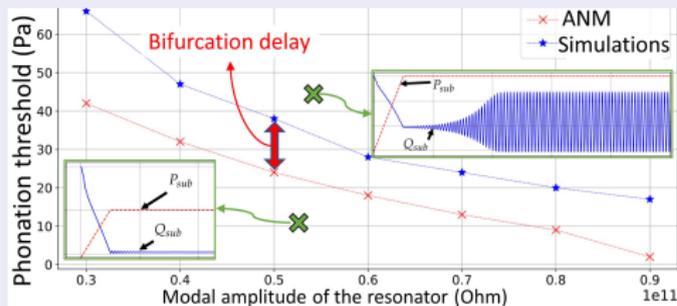
for a sufficient degree of asymmetry



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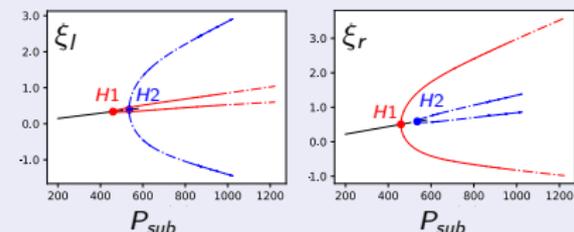
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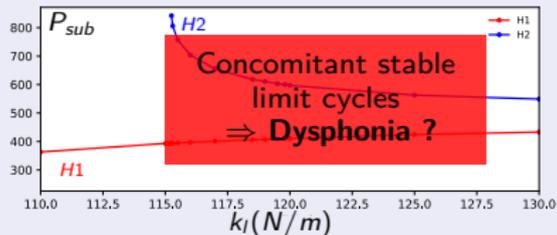
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## General conclusion and perspectives

### Models and methods (*in progress*)

- 1 Nonlinear materials and damping [Matignon, Roze]
- 2 Various families of power-balanced numerical schemes [Muller, Roze]
- 3 Automatic code generators for real-time applications [Falaize, Muller]
- 4 Passive Finite-Time Control Design [d'Andréa-Novel, Lebrun, Roze, Wijnand]
- 5 Digital Passive Controller for hardware applications [Lebrun]
- 6 Voice physical-based synthesis [Silva, Wetzel]
- 7 Regime analysis of self-oscillating PHS [Silva, Terrien, Wetzel]

### Projects based on PHS

- 1 Audio/Acoustics Virtual Factory
- 2 Augmented/Hybrid Musical Instruments with hardware development
- 3 Reprogrammed transducers (ideal HIFI loudspeaker, acoustic absorber, etc)

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– The end –

**Thank you for your attention**